

# Performance of SIM-MDPSK FSO Systems with Hardware Imperfections

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**Abstract**—This paper studies the error performance of free-space optical (FSO) systems, employing subcarrier intensity modulation (SIM) with M-ary differential phase-shift keying (MDPSK). Novel analytical expressions for the symbol error probability (SEP) are derived, based on the Fourier series approach. The irradiance fluctuations of the received optical signal are modeled by considering both Gamma-Gamma atmospheric turbulence and pointing errors. In addition, hardware imperfections of DPSK demodulator, as the phase noise of local oscillator at the receiver, are taken into account. It is illustrated that the phase noise significantly degrades the system performance, especially when the optical signal transmission is impaired by weak atmospheric turbulence and weak pointing errors effect. Furthermore, the phase noise results in an unrecoverable error-rate floor, which is an important limiting factor for SIM-DPSK FSO systems.

**Index Terms**—Atmospheric turbulence, free-space optics (FSO), Gamma-Gamma distribution, differential phase-shift keying (DPSK), phase noise, subcarrier intensity modulation (SIM), symbol error probability (SEP).

## I. INTRODUCTION

**B**ESIDES the main advantages, as high data rate, wide bandwidth and license-free transmission, free-space optical (FSO) systems are also characterized by low-power and low-cost transmission, as well as easy and simple installation. Intensity-modulation/direct detection (IM/DD) with on-off keying (OOK) is usually employed in commercial FSO systems. However, in order to improve the system performance subcarrier intensity modulation (SIM) was proposed, where the radio-frequency (RF) subcarrier signal is firstly premodulated by the data sequence bearing information, and then it is used to modulate the intensity of the laser source [1]–[4].

Several well-known modulation techniques from the field of RF communications, were used to modulate a subcarrier signal in FSO systems. The SIM based FSO system employing quadrature amplitude modulation (QAM) were analysed in [5]–[9], while SIM with M-ary phase-shift keying (MPSK) was investigated in [4], [9]–[14]. Furthermore, practical wireless communication systems also employ differential phase-shift keying (DPSK), which does not require the carrier phase estimation at the receiver. The performance of FSO systems

with coherent detection and binary DPSK (BDPSK) was analyzed in [3], [15]–[19], while the case of SIM-BDPSK was investigated in [20]–[26]. Furthermore, in order to increase the capacity (or the system throughput), an FSO system based on SIM and higher-order DPSK modulation was also proposed and analyzed in [12], [13], [27]. Specifically, an expression in integral form for the bit error rate (BER) was presented in [12], [13], while [27] compares the performance of different modulation formats, including BDPSK and quaternary DPSK (QDPSK), when space diversity is used at the reception.

The FSO system performance can be notably degraded due to the hardware imperfections. For example, the effects of the imperfect reference carrier signal phase recovery on error performance of SIM-MPSK FSO systems were examined in [28], considering weak atmospheric turbulence modeled by log-normal distribution. The effect of noisy reference signal extraction on error rate degradation of coherent BPSK FSO system in strong turbulence conditions was examined in [29]. Although the DPSK receiver does not require a carrier phase estimation, the hardware imperfections of the DPSK demodulator can seriously degrade the system performance. After optical-to-electrical signal conversion in SIM-DPSK receiver, it is necessary to down-convert the received DPSK signal. In other words, a local oscillator, used in DPSK receiver for down-conversion, generates signal, which is not ideal, in the sense that phase of this signal is a random process fluctuating over time. These fluctuations, which are in the same frequency band with the useful signal, have the influence on the detection process. This undesired phase is known as a *phase noise* [30], [31].

Scanning the open literature, to the best of the authors' knowledge, the effect of hardware imperfections as the phase noise on the performance of the FSO system employing SIM-MDPSK, has not been investigated so far. In this paper, we derive novel analytical expressions for the symbol error probability (SEP) of the SIM-MDPSK based FSO system, when hardware imperfections are considered, using the Fourier series method (FSM) [32]–[35]. The impact of hardware imperfections is represented through the phase noise, which is modeled by the Tikhonov distribution [36]–[38], and is generated by the local oscillator of DPSK demodulator [30], [31], [36], [39], [40]. The intensity fluctuations of the received optical signal are assumed to originate from the combined effect of the Gamma-Gamma atmospheric turbulence and the pointing errors [16], [41]–[45]. The derived SEP expression is given in the convergent series form, whose upper bound for the truncation error is estimated. Furthermore, the derived

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expressions are simplified, when the pointing errors effect can be neglected. Finally, numerical results are presented and validated through Monte Carlo simulations.

The rest of the paper is organized as follows. Section II describes the system and channel model, while the error analysis is provided in Section III. Numerical results with discussion are presented in Section IV and some concluding remarks are given in Section V.

## II. SYSTEM AND CHANNEL MODEL

The block diagram of the SIM-MDPSK FSO system is presented in Fig. 1. The information data are differently encoded and PSK is applied in an RF domain [46, p. 333]. DC bias is added to avoid clipping and distortion, and resulting signal modulates the laser output, by using SIM. The radiated optical power is given by

$$P(t) = P_t(1 + ms(t)), \quad (1)$$

where  $P_t$  represents the transmitted optical power and  $m$  denotes the modulation index ( $0 < m \leq 1$ ). The optical transmission via free space is influenced by atmospheric turbulence and pointing errors. At the receiver, direct detection is performed, DC bias is removed and an optical-to-electrical conversion is applied via a PIN photodetector. The electrical signal at the input of DPSK demodulator is expressed as

$$r_e(t) = I\eta P_t ms(t) + n(t), \quad (2)$$

where  $I$  is a random variable (RV), which follows Gamma-Gamma distribution and represents atmospheric turbulence and pointing errors,  $\eta$  denotes an optical-to-electrical conversion coefficient and  $n(t)$  is an additive white Gaussian noise (AWGN), with zero mean and variance,  $\sigma_n^2$ . Finally, the electrical signal,  $r_e(t)$ , is recovered by the DPSK demodulator, presented in Fig. 1, assuming that hardware imperfections exist.

### A. Modeling the combined effect of atmospheric turbulence and pointing errors

The well-known Gamma-Gamma distribution is used for describing the effect of atmospheric turbulence [41], while the pointing errors effect is described by the distribution which assumes the radial displacement of laser beam at receiver experiences Rayleigh distribution, with the jitter variance  $\sigma_s^2$  [42, (11)].

Based on (2), the instantaneous SNR is defined as  $\gamma = I^2 \eta^2 P_t^2 m^2 / (2\sigma_n^2)$ . The probability density function (PDF) of  $\gamma$  is [5]

$$f_\gamma(\gamma) = \frac{\xi^2}{2\Gamma(\alpha)\Gamma(\beta)\gamma} G_{1,3}^{3,0} \left( \alpha\beta\kappa \sqrt{\frac{\gamma}{\mu}} \middle| \begin{matrix} \xi^2+1 \\ \xi^2, \alpha, \beta \end{matrix} \right), \quad (3)$$

where  $G_{p,q}^{m,n}(\cdot)$  is the Meijer's  $G$ -function [47, (9.301)], and  $\mu$  represents the average electrical SNR per symbol. The relation between  $\mu$  and the average electrical SNR per bit,  $\mu_b$ , is  $\mu = \mu_b \log_2 M$ . The average electrical SNR per bit is defined as  $\mu_b = \eta^2 P_t^2 m^2 \kappa^2 A_0^2 I_1^2 / (2\sigma_n^2)$ , with  $\kappa = \xi^2 / (\xi^2 + 1)$  [5]. The atmospheric turbulence parameters are denoted by  $\alpha$  and  $\beta$ , while  $\xi$  and  $A_0$  represent the pointing errors parameters.

Assuming Gaussian plane wave propagation and zero inner scale, the parameters  $\alpha$  and  $\beta$  are defined as  $\alpha = (\exp[0.49\sigma_R^2(1 + 1.11\sigma_R^{12/5})^{-7/6}] - 1)^{-1}$  and  $\beta = (\exp[0.51\sigma_R^2(1 + 0.69\sigma_R^{12/5})^{-5/6}] - 1)^{-1}$  [1], [41], with the Rytov variance  $\sigma_R^2 = 1.23C_n^2 k^{7/6} L^{11/6}$ . The wave-number is  $k = 2\pi/\lambda$  with the wavelength  $\lambda$ ,  $L$  is the propagation distance, and the refractive index is denoted by  $C_n^2$ .

The pointing error represents the misalignment between the transmitter laser and the receiver photodetector. The parameter  $\xi$  is defined as the ratio between the equivalent beam radius at the receiver,  $w_{Leq}$ , and the pointing error (jitter) standard deviation at the receiver as  $\xi = w_{Leq}/(2\sigma_s)$ . The parameter,  $w_{Leq}$ , depends on the beam radius at distance  $L$ ,  $w_L$ , as  $w_{Leq}^2 = w_L^2 \sqrt{\pi} \operatorname{erf}(v) / (2v \exp(-v^2))$ ,  $v = \sqrt{\pi} a / (\sqrt{2} w_L)$  [42], where  $a$  is the radius of a circular detector aperture,  $\operatorname{erf}(\cdot)$  is the error function [47, (8.250.1)], and  $A_0 = [\operatorname{erf}(v)]^2$ . Next, the parameter  $w_L$  is related with the beam radius at the waist,  $w_0$ , and the radius of curvature,  $F_0$ , by  $w_L = w_0((\Theta_o + \Lambda_o)(1 + 1.63\sigma_R^{12/5}\Lambda_1))^{1/2}$ , where  $\Theta_o = 1 - L/F_0$ ,  $\Lambda_o = 2L/(kw_0^2)$ ,  $\Lambda_1 = \Lambda_o/(\Theta_o^2 + \Lambda_o^2)$  [44].

### B. Phase noise

After signal conversion from optical-to-electrical domain, classical signal detection is performed in electrical domain. During the process of down-conversion, electrical signal is multiplied by local oscillator output signal. The phase of the local oscillator signal (also known as a phase noise) is a random process fluctuating over time. Frequently local oscillator is embedded in frequency syntetyzator contained phase locked loop (PLL). The phase noise generated by PLL is well known to have a Tikhonov PDF [37, Ch. 2], [38]. Hence, the phase noise,  $\varphi$ , of the local oscillator is assumed to be a RV which follows Tikhonov PDF given by

$$f_\varphi(\varphi) = \frac{\exp(b \cos(\varphi))}{2\pi I_0(b)}, \quad |\varphi| \leq \pi, \quad (4)$$

where  $I_n(\cdot)$  is the  $n$ th order modified Bessel function of the first kind [47, (8.431)],  $b = 1/\sigma_\varphi^2$ , and  $\sigma_\varphi^2$  is the variance of the phase noise.

Here, we use the Fourier series expansion of Tikhonov PDF, because Fourier series form is tractable for integration that will be necessary in mathematical derivations of SEP. We start with the Fourier expansion [48, (9.6.34)]

$$e^{b \cos \varphi} = I_0(b) + 2 \sum_{n=1}^{\infty} I_n(b) \cos(n\varphi), \quad |\varphi| \leq \pi, \quad (5)$$

for a fixed  $b > 0$ .

Based on the expansion in (5), it is clear that Tikhonov PDF given by (4), can be expressed in the Fourier series as

$$f_\varphi(\varphi) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} c_n \cos(n\varphi), \quad |\varphi| \leq \pi, \quad (6)$$

where

$$c_n = \frac{I_n(b)}{\pi I_0(b)}. \quad (7)$$

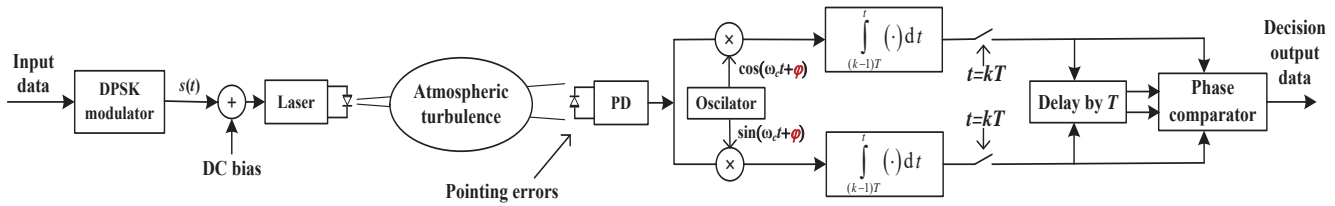


Fig. 1. Block diagram of a SIM-MDPSK FSO system

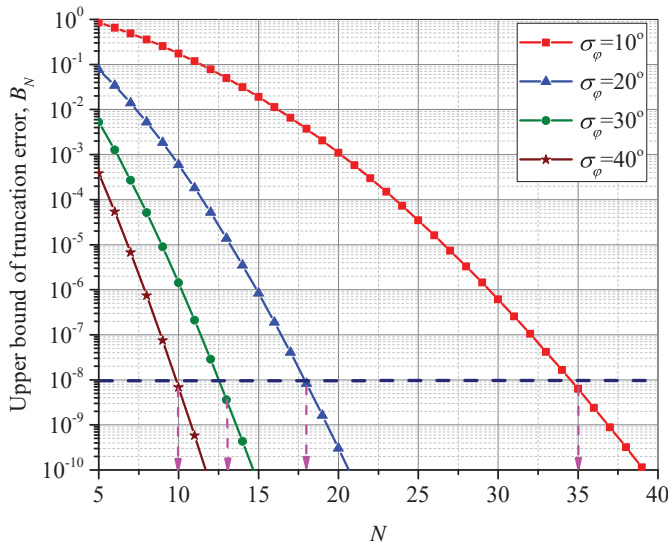


Fig. 2. Upper bound of truncation errors for  $\sigma_\varphi = 10^\circ$  (red),  $\sigma_\varphi = 20^\circ$  (blue),  $\sigma_\varphi = 30^\circ$  (green), and  $\sigma_\varphi = 40^\circ$  (brown) when  $N \leq 40$

*Proposition 1:* The series in (6) is convergent. For the truncation error

$$E_N(\varphi; b) = \sum_{n=N+1}^{\infty} c_n \cos(n\varphi), \quad |\varphi| \leq \pi, \quad (8)$$

the following estimate

$$|E_N(\varphi; b)| \leq E_N(0; b) \leq B_N \quad (9)$$

holds, where

$$B_N \equiv B_N(b) = \frac{1}{\pi I_0(b)} \left( I_{N+1}(b) + \int_{N+1}^{\infty} I_\nu(b) d\nu \right). \quad (10)$$

*Proof:* See Appendix A.

In Fig. 2, we present the bounds  $B_N$  of the truncation errors for  $N \leq 40$  and different values of  $\sigma_\varphi$ . If we take a threshold for the errors, e.g.,  $\varepsilon = 10^{-8}$  (black line in Fig. 2), so that  $B_N < \varepsilon$ , we see that the corresponding number of terms should be  $N = 35, 18, 13$  and  $10$  for  $\sigma_\varphi = 10^\circ, 20^\circ, 30^\circ$  and  $40^\circ$ , respectively.

### III. ERROR PERFORMANCE

Since the decisions of the DPSK receiver are taken based on the composite phase difference between signals received

during two consecutive symbol intervals, the decision variable of differential detector can be written as

$$\lambda' = [\psi_{k+1} - \psi_k] \bmod 2\pi, \quad (11)$$

where  $\psi_{k+1}$  and  $\psi_k$  are the composite phase of consecutive received signals, bearing the information at the  $(k+1)$ -th and the  $k$ -th interval, respectively. The local oscillator imperfections are represented through the phase noise  $\varphi_{k+1}$  and  $\varphi_k$  at the  $(k+1)$ -th and the  $k$ -th intervals, respectively. Then, the decision variable of the differential detector is

$$\begin{aligned} \lambda &= [(\psi_{k+1} - \varphi_{k+1}) - (\psi_k - \varphi_k)] \bmod 2\pi \\ &= [(\psi_{k+1} - \psi_k) - (\varphi_{k+1} - \varphi_k)] \bmod 2\pi. \end{aligned} \quad (12)$$

The term,  $(\psi_{k+1} - \psi_k)$ , represents the difference of the composite phases, while,  $(\varphi_{k+1} - \varphi_k)$ , denotes the impact of the phase noise.

On the contrary to the situation at the transmitter, where the phase of RF carrier is constant, the composite phase of total received signal is a RV. The PDF of the resulting phase,  $\psi$ , of received signal in a signaling interval is presented in the Fourier series form as [32]–[34]

$$f_\psi(\psi) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} b_n \cos(n\psi), \quad (13)$$

where  $b_n$  represents the Fourier coefficient for the FSO channel influenced by the Gamma-Gamma atmospheric turbulence and pointing errors. In order to derive the Fourier coefficient for the considered scenario, the PDF of the received signal composite phase is written as

$$f_\psi(\psi) = \int_0^{\infty} f(\psi|\gamma) f_\gamma(\gamma) d\gamma \quad (14)$$

where  $f_\gamma(\gamma)$  is the PDF of the instantaneous SNR given in (3). The conditional PDF is defined through a Fourier series form of the received signal composite phase due to additive noise as [32]–[34]

$$f(\psi|\gamma) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} a_n(\gamma) \cos(n\psi), \quad (15)$$

where  $a_n(\gamma)$  denotes the Fourier coefficient for AWGN channel defined as [34]

$$a_n(\gamma) = \frac{\Gamma\left(\frac{n}{2}+1\right)}{n!\pi} \gamma^{\frac{n}{2}} \exp(-\gamma) {}_1F_1\left(\frac{n}{2}+1; n+1; \gamma\right), \quad (16)$$

where  ${}_1F_1(\cdot; \cdot; \cdot)$  is the confluent hypergeometric function [47, (9.21)].

*Proposition 2:* After substituting (3), (15) and (16) into (14), the PDF of the phase  $\psi$  is given as

$$f_{\psi}(\psi) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} \frac{n2^{\alpha+\beta-4}\xi^2}{\pi^2\Gamma(\alpha)\Gamma(\beta)} \cos(n\psi) \times G_{3,6}^{6,1} \left( \frac{\alpha^2\beta^2\kappa^2}{16\mu} \left| \begin{matrix} 1-\frac{n}{2}, 1+\frac{n}{2}, \frac{\xi^2+2}{2} \\ \frac{\xi^2}{2}, \frac{\alpha}{2}, \frac{\alpha+1}{2}, \frac{\beta}{2}, \frac{\beta+1}{2}, 0 \end{matrix} \right. \right). \quad (17)$$

*Proof:* See Appendix B.

Based on (13) and (17), the Fourier coefficient for FSO channel influenced by Gamma-Gamma atmospheric turbulence and pointing errors is determined as

$$b_n = \frac{n2^{\alpha+\beta-4}\xi^2}{\pi^2\Gamma(\alpha)\Gamma(\beta)} \times G_{3,6}^{6,1} \left( \frac{\alpha^2\beta^2\kappa^2}{16\mu} \left| \begin{matrix} 1-\frac{n}{2}, 1+\frac{n}{2}, \frac{\xi^2+2}{2} \\ \frac{\xi^2}{2}, \frac{\alpha}{2}, \frac{\alpha+1}{2}, \frac{\beta}{2}, \frac{\beta+1}{2}, 0 \end{matrix} \right. \right). \quad (18)$$

When the considered scenario assumes the pointing errors to be very small, it can be neglected ( $\xi \rightarrow \infty$ ). In this case, the optical link suffers only from atmospheric turbulence, and the Fourier coefficient can be found by taking the limit of (18) for  $\xi \rightarrow \infty$ . After applying [49, (07.34.25.0007.01), (07.34.25.0006.01) and (06.05.16.0002.01)], the Fourier coefficient can be derived as

$$b_n^{GG} = \lim_{\xi \rightarrow \infty} b_n = \frac{n2^{\alpha+\beta-3}}{\pi^2\Gamma(\alpha)\Gamma(\beta)} \times G_{2,5}^{5,1} \left( \frac{\alpha^2\beta^2}{16\mu} \left| \begin{matrix} 1-\frac{n}{2}, 1+\frac{n}{2} \\ \frac{\alpha}{2}, \frac{\alpha+1}{2}, \frac{\beta}{2}, \frac{\beta+1}{2}, 0 \end{matrix} \right. \right). \quad (19)$$

For further analysis, it is required to find the PDF of the decision variable  $\lambda$ , defined in (12). Firstly, we will introduce the following rule related to the PDFs presented in the Fourier series form.

*Proposition 3:* If the variables  $x_1$  and  $x_2$  are RVs with the PDFs given in the Fourier series form, with coefficients  $z_{1n}$  and  $z_{2n}$ , respectively, as

$$f_{x_1}(x) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} z_{1n} \cos(nx), \quad |x| \leq \pi, \quad (20)$$

$$f_{x_2}(x) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} z_{2n} \cos(nx), \quad |x| \leq \pi,$$

then, the PDF of  $y = [x_1 - x_2] \bmod 2\pi$ , is

$$f_y(y) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} \pi z_{1n} z_{2n} \cos(ny), \quad |y| \leq \pi. \quad (21)$$

*Proof:* The proof can be found in [36], [50], [51].

#### A. Error analysis without considering hardware imperfections

If no hardware imperfections are assumed, the decision variable  $\lambda'$  is defined in (11). Based on Proposition 3, after replacing  $x_1$  and  $x_2$  with  $\psi_{k+1}$  and  $\psi_k$ , respectively, and both  $z_{1n}$  and  $z_{2n}$  with  $b_n$ , the PDF of  $\lambda'$  can be easily obtained as

$$f_{\lambda'}(\lambda') = \frac{1}{2\pi} + \sum_{n=1}^{\infty} \pi b_n^2 \cos(n\lambda'), \quad |\lambda'| \leq \pi. \quad (22)$$

The detection is performed in the manner to find the closest possible transmitted phase compared with received composite phase  $\lambda'$ . The probability of wrong symbol detection is given by

$$P_s = 1 - \int_{-\pi/M}^{\pi/M} f_{\lambda'}(\lambda') d\lambda'. \quad (23)$$

By substituting (22) into (23), the average SEP can be found as

$$P_s = 1 - \frac{1}{M} - \sum_{n=1}^{\infty} \frac{2\pi b_n^2}{n} \sin\left(\frac{n\pi}{M}\right). \quad (24)$$

In [13], an expression for the average BER was derived in integral form, assuming that the intensity fluctuations of the optical signal are modeled by the log-normal and Gamma-Gamma distributions. In the region of high average electrical SNR values, the bit error probability could be approximated by  $BER \approx P_s / \log_2 M$  [40, p. 271]. By using this approximation and SEP in (24) with the Fourier coefficient of (19), the numerical results from [13, Fig. 2] can be obtained.

#### B. Error analysis in the presence of phase noise

In the presence of the phase noise, the decision variable  $\lambda$  is defined as in (12). The PDF of the variable,  $\delta = \varphi_{k+1} - \varphi_k$ , can be found by utilization of Proposition 3. Since the Tikhonov PDF of the phase noise is given in the Fourier series form by (6), the PDF of the variable  $\delta$  is found as

$$f_{\delta}(\delta) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} \pi c_n^2 \cos(n\delta), \quad |\delta| \leq \pi, \quad (25)$$

with the Fourier coefficient  $c_n$  previously defined by (7).

Taking into consideration that the variables  $\psi$  and  $\varphi$  are statistically independent, based on (22) and (25), and Proposition 3, the PDF of  $\lambda$  is

$$f_{\lambda}(\lambda) = \frac{1}{2\pi} + \sum_{n=1}^{\infty} \pi^3 b_n^2 c_n^2 \cos(n\lambda), \quad |\lambda| \leq \pi. \quad (26)$$

When the Gamma-Gamma atmospheric turbulence, pointing errors and phase noise are assumed, the average SEP of the SIM-MDPSK FSO system can be written as

$$P_s = 1 - \int_{-\pi/M}^{\pi/M} f_{\lambda}(\lambda) d\lambda = 1 - \frac{1}{M} - \sum_{n=1}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} \sin\left(\frac{n\pi}{M}\right), \quad (27)$$

where the Fourier coefficients  $b_n$  and  $c_n$  are previously defined in (18) and (7), respectively.

*Proposition 4:* The series in (27) is convergent and the following estimate

$$\left| P_s - 1 + \frac{1}{M} + \sum_{n=1}^N \frac{2\pi^3 b_n^2 c_n^2}{n} \sin\left(\frac{n\pi}{M}\right) \right| \leq E_N^{\text{SEP}} \quad (28)$$



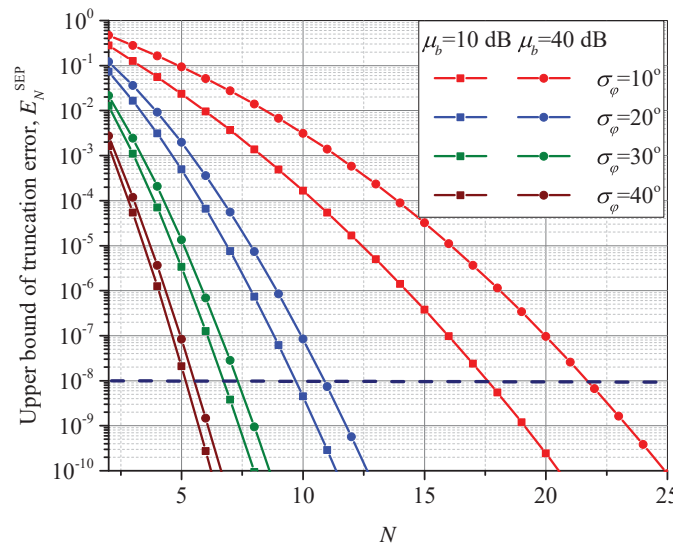


Fig. 3. Upper bound of truncation errors for  $\sigma_\varphi = 10^\circ$  (red),  $\sigma_\varphi = 20^\circ$  (blue),  $\sigma_\varphi = 30^\circ$  (green), and  $\sigma_\varphi = 40^\circ$  (brown) when  $N \leq 25$

holds, with the bound of truncation error

$$E_N^{\text{SEP}} = \frac{2\pi b_{N+1}^2}{I_0(b)^2} \left( \frac{I_{N+1}(b)^2}{N+1} + \int_{N+1}^{\infty} \frac{I_\nu(b)^2}{\nu} d\nu \right). \quad (29)$$

*Proof:* See Appendix C.

This truncation error is illustrated in Fig. 3. To achieve the given truncation error, the higher number of terms in summation is required if the standard deviation is lower. In order to achieve truncation error less than  $10^{-8}$ , for  $\mu = 10$  dB the required number of terms in summation is  $N = 18, 10, 7$  and  $6$ , when  $\sigma_\varphi = 10^\circ, 20^\circ, 30^\circ$  and  $40^\circ$ , respectively. In addition, the convergence rate decreases with increasing the electrical SNR. In other words, the proposed series expression converges better in low electrical SNR regime compared to high electrical SNR regime.

As it will be shown in the next Section, the existence of the phase noise results in the unrecoverable error-rate floor, which is a meaningful limiting factor in SIM-DPSK based FSO systems. This error-rate floor represents the constant value of the average SEP, which occurs at the high average electrical SNR. With a further increase in the transmitted optical power, the improvement of the SEP performance will not be achieved.

*Proposition 5:* The unrecoverable error-rate floor can be expressed as

$$P_s^{\text{floor}} = 1 - \frac{1}{M} - \sum_{n=1}^{\infty} \frac{2\pi c_n^2}{n} \sin\left(\frac{n\pi}{M}\right). \quad (30)$$

*Proof:* See Appendix D.

It can be noticed that the SEP floor is independent on the FSO channel state (atmospheric turbulence and pointing errors). On the other hand, the value of the SEP floor depends on the phase noise standard deviation and order of DPSK modulation, as it will be presented in the next Section.

#### IV. NUMERICAL RESULTS AND DISCUSSION

Based on derived expressions for the average SEP, numerical results are obtained and validated by Monte Carlo

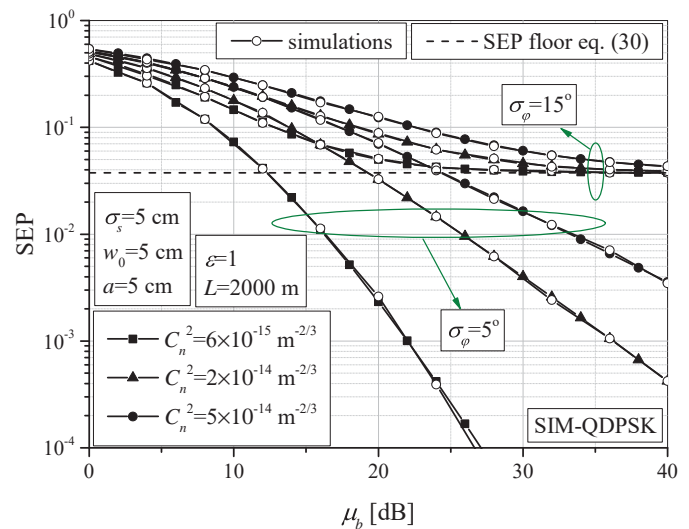


Fig. 4. SIM-QDPSK SEP versus average electrical SNR for different values of the phase noise standard deviation in various atmospheric turbulence conditions

simulations. Monte Carlo simulations have been performed using MATLAB<sup>®</sup> software package. Since intensity fluctuations originate from both atmospheric turbulence and pointing errors, the resulting optical signal intensity,  $I$ , is obtained as a product of two different RVs, i.e.,  $I = I_a \times I_p$ . The intensity fluctuations,  $I_a$ , due to atmospheric turbulence are modeled by Gamma-Gamma distribution. The corresponding RV,  $I_a$ , is generated as a product of two independent Gamma-distributed RVs with shaping parameters  $\alpha$  and  $\beta$ . Command for generating Gamma-distributed RV is built-in into MATLAB<sup>®</sup>. The RVs relating to the pointing errors,  $I_p$ , are generated based on [42, (9)], employing built-in command for generating Rayleigh RV. The Tikhonov-distributed samples of phase noise are generated using the modified acceptance/rejection method, explained in [52, p. 382]. Modulation and demodulation is simulated based on [46, p. 333-335]. The average SEP values are estimated using  $10^7$  transmitted symbols.

In order to obtain the numerical results, the atmospheric turbulence strength is determined by the refractive index structure parameter as:  $C_n^2 = 6 \times 10^{-15} \text{ m}^{-2/3}$  for weak,  $C_n^2 = 2 \times 10^{-14} \text{ m}^{-2/3}$  for moderate and  $C_n^2 = 5 \times 10^{-14} \text{ m}^{-2/3}$  for strong turbulence conditions. The impact of the phase noise is specified by the phase noise standard deviation.

The average SEP dependence on the average electrical SNR of the FSO system employing SIM-QDPSK is presented in Fig. 4, assuming different atmospheric turbulence conditions and phase noise standard deviation  $\sigma_\varphi = 5^\circ$  or  $\sigma_\varphi = 15^\circ$ . Lower values of the phase noise standard deviation correspond to the weaker phase noise and better system performance. Furthermore, the impact of the atmospheric turbulence conditions is stronger when the value of  $\sigma_\varphi$  is lower. On the other hand, when the effect of phase noise is very strong, the atmospheric turbulence conditions has minor influence on the SEP performance. In addition, the existence of the unrecoverable error-rate floor is noticed in Fig. 4, meaning that the DPSK hardware imperfections presented through phase

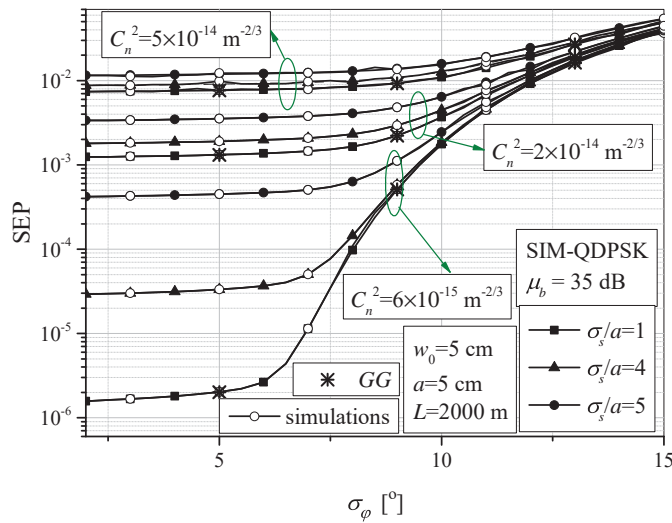


Fig. 5. SIM-QDPSK SEP versus the phase noise standard deviation for different values of the normalized jitter standard deviation, in various atmospheric turbulence conditions

noise are an important limiting factor for SIM-DPSK systems.

This SEP floor appears at lower values of average electrical SNR in weak atmospheric turbulence, as well as when the value of  $\sigma_\varphi$  is greater (stronger impact of the phase noise). The SEP floor results based on (30) for  $\sigma_\varphi = 5^\circ$  are not visible in Fig. 4 due to very low value. It can be concluded that the SEP floor is not dependent on atmospheric turbulence conditions, which is in agreement with mathematical derivation (see (41) and (30)).

Fig. 5 presents the SIM-QDPSK SEP dependence on the phase noise standard deviation for different values of the normalized jitter standard deviation, in various atmospheric turbulence conditions. It can be observed that lower values of the normalized jitter standard deviation reflects in better system performance. It means that the positioning of the FSO apertures is better and the pointing errors effect is weaker. Also, the pointing error effect is stronger in weak compared to moderate and strong atmospheric turbulence. When the optical signal transmission suffers from very strong atmospheric turbulence, the pointing errors effect has less impact on the SEP performance.

In addition, the results for the FSO system when the pointing errors effect is neglected, obtained by using (27) and (19), are also presented. These results are in agreement with those when  $\sigma_s/a = 1$ . Hence, very low values of the normalized jitter standard deviation means that the pointing errors effect is very weak and can be neglected.

When the DPSK demodulator hardware imperfections are dominant, and the phase noise is quite strong, the value of  $\sigma_\varphi$  is large. In that case, the FSO channel state (atmospheric turbulence and pointing errors) does not play a major role in the SEP performance. When  $\sigma_\varphi \rightarrow 0$ , the impact of the phase noise is very weak and can be neglected. For these phase noise standard deviation values, the SEP takes constant values, which are approximately the same as the SEP values for the FSO system without phase noise. Also, atmospheric turbulence and pointing errors have very strong impact on the

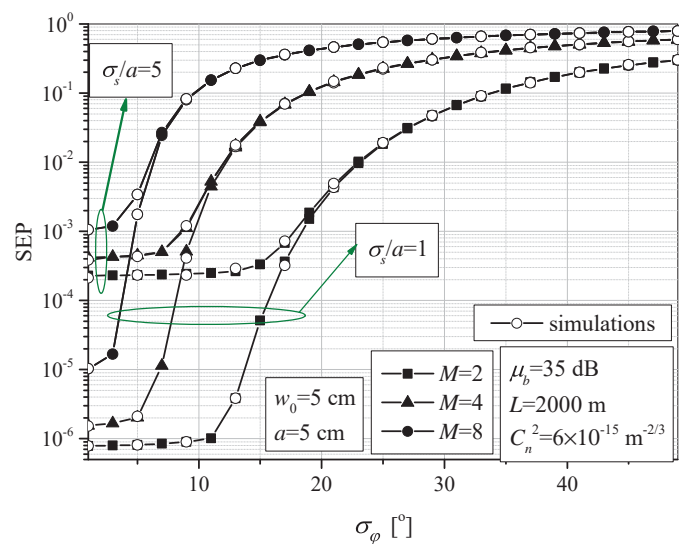


Fig. 6. SIM-MDPSK SEP versus the phase noise standard deviation for different values of the normalized jitter standard deviations

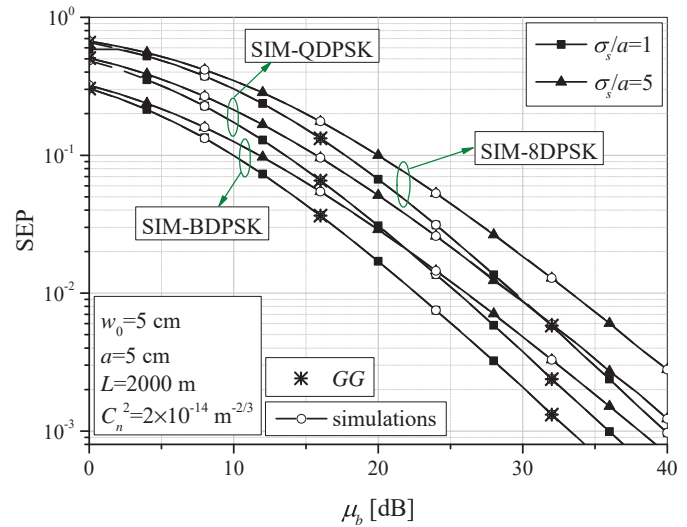


Fig. 7. SIM-MDPSK SEP versus average electrical SNR of the FSO system without hardware imperfections

SEP performance, when  $\sigma_\varphi$  is low.

Fig. 6 represents the SIM-MDPSK SEP dependence on the phase noise standard deviation. The impact of the phase noise on SEP is stronger when higher order SIM-MDPSK is employed. For example, for  $\sigma_s/a = 1$ , in the case of  $M = 2$ , the SEP is independent on phase noise up to  $\sigma_\varphi = 10^\circ$ , while for  $M = 8$ , SEP drastically increases even starting from  $\sigma_\varphi = 2^\circ$ . In addition, the weaker the pointing errors, the stronger is the effect of phase noise on SEP. It can be observed that the effect of DPSK order has minor influence on the SEP performance when the impact of the phase noise is very strong.

The SEP dependence on the average electrical SNR of the FSO system without hardware imperfections is presented in Fig. 7. The results are obtained based on (24) with the Fourier coefficient in (18), or in (19) when the pointing errors

are neglected. Different DPSK formats are observed: SIM employing BDPSK, QDPSK and 8DPSK. As it is expected, FSO system based on SIM-DPSK with higher modulation format has worse SEP performance, but the larger amount of information can be transmitted. Also, consistent with previous conclusions, greater value of the normalized jitter standard deviations means worse system performance due to stronger pointing errors. Agreement of the results based on (18) for  $\sigma_s/a = 1$  and (19) is noticed, meaning that very low jitter standard deviation leads to weak pointing errors.

## V. CONCLUSION

We have derived novel analytical expressions for the average SEP of FSO system employing SIM-MDPSK. The irradiance fluctuations at the received signal originate from the Gamma-Gamma atmospheric turbulence and pointing errors. Based on derived SEP expressions, numerical results have been presented and confirmed by Monte Carlo simulations.

From the illustrated results, we have found that the hardware imperfections result in the significant deterioration of the FSO system performance. The phase noise is dominant system factor, which causes the SEP performance damaging, especially when optical signal transmission is influenced by favorable conditions (weak atmospheric turbulence and weak pointing errors effect). Similarly, when the impact of the phase noise is very strong, atmospheric turbulence and pointing errors effect has minor effect on the system performance. Furthermore, the SIM based FSO system with higher DPSK format is more sensitive to the existence of the phase noise. Further, the existence of the phase noise leads to the unrecoverable SEP floor, being meaningful limiting factor for SIM-DPSK systems. It is observed that the SEP floor is not dependent on the FSO channel state, but it is highly dependent on the phase noise standard deviation and the DPSK modulation order.

## APPENDIX A PROOF OF PROPOSITION 1

The series in (6) is a uniformly convergent series, because the numerical series with positive terms,

$$\sum_{n=1}^{\infty} I_n(b), \quad (31)$$

is convergent, which can be proved using the inequality [53]

$$\left(1 + \frac{\nu}{b}\right) I_{\nu+1}(b) < I_{\nu}(b) \quad (\nu \geq -1, b > 0). \quad (32)$$

Namely, the series (31) is convergent if for a fixed  $m = \lceil b \rceil$  ( $\lceil x \rceil$  denotes the smallest integer greater than or equal to  $x$ ) the series  $\sum_{n=m}^{\infty} I_n(b)$  converges. According to (32), for  $n > m$  we have

$$I_n(b) < \frac{I_{n-1}(b)}{1 + (n-1)/b} < \frac{I_{n-1}(b)}{2} < \dots < \frac{I_m(b)}{2^{n-m}}, \quad (33)$$

so that

$$\sum_{n=m}^{\infty} I_n(b) < I_m(b) \sum_{n=m}^{\infty} \frac{1}{2^{n-m}} = 2I_m(b),$$

wherefrom we conclude that the series  $\sum_{n=m}^{\infty} I_n(b)$  and (31) are convergent. The sum of (31) is  $S = \frac{1}{2}(e^b - I_0(b))$  [54, p. 254].

According to Cauchy's integral test (cf. [55, p. 120] or [56, p. 159]), for the numerical series (31) we can give the following estimates for the remainder term

$$\sum_{n=N+1}^{\infty} I_n(b) \leq I_{N+1}(b) + \int_{N+1}^{\infty} I_{\nu}(b) d\nu, \quad (34)$$

where  $\nu \mapsto I_{\nu}(a)$  is a decreasing positive continuous function on  $(0, \infty)$  [53].

Thus, for the truncation error  $E_N(\varphi; b)$  given by (8) we obtain  $|E_N(\varphi; b)| \leq E_N(0; b) = (\pi I_0(b))^{-1} \sum_{n=N+1}^{\infty} I_n(b)$ , i.e., (9), where  $B_N$  is given by (10), because of (34).

## APPENDIX B PROOF OF PROPOSITION 2

After substituting (3), (15) and (16) into (14), the PDF of the phase  $\psi$  is re-written as

$$\begin{aligned} f_{\psi}(\psi) &= \frac{1}{2\pi} + \sum_{n=1}^{\infty} \frac{\Gamma(\frac{n}{2} + 1)}{n! \pi} \frac{\xi^2}{2\Gamma(\alpha)\Gamma(\beta)} \cos(n\psi) \\ &\times \int_0^{\infty} \gamma^{\frac{n}{2}-1} \exp(-\gamma) {}_1F_1\left(\frac{n}{2} + 1; n + 1; \gamma\right) \\ &\times G_{1,3}^{3,0}\left(\alpha\beta\kappa\sqrt{\frac{\gamma}{\mu}} \middle| \begin{matrix} \xi^2+1 \\ \xi^2, \alpha, \beta \end{matrix}\right) d\gamma. \end{aligned} \quad (35)$$

Based on [49, (07.20.26.0015.01)], the product of exponential and confluent hypergeometric function is presented in terms of the Meijer's  $G$ -function as

$$\exp(-\gamma) {}_1F_1\left(\frac{n}{2} + 1; n + 1; \gamma\right) = \frac{\Gamma(n+1)}{\Gamma(\frac{n}{2})} G_{1,2}^{1,1}\left(\gamma \middle| \begin{matrix} 1-\frac{n}{2} \\ 0, -n \end{matrix}\right). \quad (36)$$

After substituting (36) into (35) and applying [49, (06.05.16.0002.01) and (06.05.03.0001.01)], the PDF of the phase  $\psi$  is

$$\begin{aligned} f_{\psi}(\psi) &= \frac{1}{2\pi} + \sum_{n=1}^{\infty} \frac{n\xi^2}{4\pi\Gamma(\alpha)\Gamma(\beta)} \cos(n\psi) \\ &\times \int_0^{\infty} \gamma^{\frac{n}{2}-1} G_{1,2}^{1,1}\left(\gamma \middle| \begin{matrix} 1-\frac{n}{2} \\ 0, -n \end{matrix}\right) G_{1,3}^{3,0}\left(\alpha\beta\kappa\sqrt{\frac{\gamma}{\mu}} \middle| \begin{matrix} \xi^2+1 \\ \xi^2, \alpha, \beta \end{matrix}\right) d\gamma. \end{aligned} \quad (37)$$

The integral in (37) can be evaluated in closed-form by using [49, (07.34.21.0013.01)], so the PDF of the phase  $\psi$  is derived as

$$\begin{aligned} f_{\psi}(\psi) &= \frac{1}{2\pi} + \sum_{n=1}^{\infty} \frac{n2^{\alpha+\beta-4}\xi^2}{\pi^2\Gamma(\alpha)\Gamma(\beta)} \cos(n\psi) \\ &\times G_{4,7}^{7,1}\left(\frac{\alpha^2\beta^2\kappa^2}{16\mu} \middle| \begin{matrix} 1-\frac{n}{2}, 1+\frac{n}{2}, \frac{\xi^2+2}{2}, \frac{\xi^2+2}{2} \\ \frac{\xi^2}{2}, \frac{\xi^2+1}{2}, \frac{\alpha}{2}, \frac{\alpha+1}{2}, \frac{\beta}{2}, \frac{\beta+1}{2}, 0 \end{matrix}\right). \end{aligned} \quad (38)$$

After the permutation of the parameters via [49, (07.34.04.0003.01) and (07.34.04.0004.01)], and the transformation of the Meijer's  $G$ -function by [49, (07.34.03.0002.01)], the final form of the PDF of the phase  $\psi$  is presented in (17).

APPENDIX C  
PROOF OF PROPOSITION 4

We observe the series in (27) given by

$$S = \sum_{n=1}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} \sin\left(\frac{n\pi}{M}\right), \quad (39)$$

for which we can prove its absolute convergence. As in APPENDIX A we use the inequalities (32) and (33) and consider the series  $\sum_{n=m}^{\infty} 2\pi^3 b_n^2 c_n^2/n$ , where  $m = \lceil b \rceil$ . Since  $b_n$  is a decreasing sequence, we can write

$$\sum_{n=m}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} < 2\pi^3 b_m^2 \sum_{n=m}^{\infty} \frac{c_n^2}{n} = \frac{2\pi b_m^2}{I_0(b)^2} \sum_{n=m}^{\infty} \frac{I_n(b)^2}{n}.$$

Now, using (33) we conclude that

$$\sum_{n=m}^{\infty} \frac{I_n(b)^2}{n} < I_m(b)^2 \sum_{n=m}^{\infty} \frac{1}{n4^{n-m}} < \frac{4I_m(b)^2}{3m},$$

i.e.,

$$\sum_{n=m}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} < \frac{8\pi b_m^2}{3m} \left( \frac{I_m(b)}{I_0(b)} \right)^2 < +\infty.$$

Thus, the series (39) is absolutely convergent, and also convergent. For its truncation error we obtain the following estimate

$$\left| \sum_{n=N+1}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} \sin\left(\frac{n\pi}{M}\right) \right| \leq \sum_{n=N+1}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} \leq \frac{2\pi b_{N+1}^2}{I_0(b)^2} \sum_{n=N+1}^{\infty} \frac{I_n(b)^2}{n}.$$

Based on Cauchy's criteria, as in APPENDIX A, it follows

$$\sum_{n=N+1}^{\infty} \frac{I_n(b)^2}{n} \leq \frac{I_{N+1}(b)^2}{N+1} + \int_{N+1}^{\infty} \frac{I_\nu(b)^2}{\nu} d\nu$$

so that we get (28), with (29).

APPENDIX D  
PROOF OF PROPOSITION 5

In order to determine the value of the SEP floor, it is necessary to take the limit of (27) for  $\mu \rightarrow \infty$ , i.e.,

$$P_s^{floor} = \lim_{\mu \rightarrow \infty} P_s = \lim_{\mu \rightarrow \infty} \left\{ 1 - \frac{1}{M} - \sum_{n=1}^{\infty} \frac{2\pi^3 b_n^2 c_n^2}{n} \sin\left(\frac{n\pi}{M}\right) \right\}. \quad (40)$$

Since the Fourier coefficient  $b_n$  is the only term in (27), which depends on the average electrical SNR, after following derivation in this Appendix, the limit of  $b_n$  for  $\mu \rightarrow \infty$  is derived as

$$b_n^{\mu \rightarrow \infty} = \lim_{\mu \rightarrow \infty} b_n = \frac{1}{\pi}. \quad (41)$$

The term  $b_n^{\mu \rightarrow \infty}$  is derived by following

$$\begin{aligned} \lim_{\mu \rightarrow \infty} b_n &= \lim_{\mu \rightarrow \infty} \frac{n2^{\alpha+\beta-4}\xi^2}{\pi^2 \Gamma(\alpha) \Gamma(\beta)} \\ &\times G_{3,6}^{6,1} \left( \frac{\alpha^2 \beta^2 \kappa^2}{16\mu} \left| \begin{matrix} 1-\frac{n}{2}, 1+\frac{n}{2}, \frac{\xi^2+2}{2} \\ \frac{\xi^2}{2}, \frac{\alpha}{2}, \frac{\alpha+1}{2}, \frac{\beta}{2}, \frac{\beta+1}{2}, 0 \end{matrix} \right. \right) \\ &= \lim_{z \rightarrow 0} \frac{n2^{\alpha+\beta-4}\xi^2}{\pi^2 \Gamma(\alpha) \Gamma(\beta)} \\ &\times G_{3,6}^{6,1} \left( z \left| \begin{matrix} 1-\frac{n}{2}, 1+\frac{n}{2}, \frac{\xi^2+2}{2} \\ \frac{\xi^2}{2}, \frac{\alpha}{2}, \frac{\alpha+1}{2}, \frac{\beta}{2}, \frac{\beta+1}{2}, 0 \end{matrix} \right. \right). \end{aligned} \quad (42)$$

The first step in finding  $\lim_{\mu \rightarrow \infty} b_n$  is applying [49, (07.34.06.0001.01)] to represent the Meijer's  $G$ -function in (42) in series form. Since  $z \rightarrow 0$ , higher order terms in the series representation of Meijer's  $G$ -function can be neglected, and  $b_n^{\mu \rightarrow \infty}$  is determined as

$$\begin{aligned} b_n^{\mu \rightarrow \infty} &= \lim_{\mu \rightarrow \infty} b_n \approx \frac{1}{2^2 \pi^2} \\ &\times \frac{2^\alpha \Gamma\left(\frac{\alpha}{2}\right) \Gamma\left(\frac{\alpha+1}{2}\right) 2^\beta \Gamma\left(\frac{\beta}{2}\right) \Gamma\left(\frac{\beta+1}{2}\right)}{\Gamma(\alpha) \Gamma(\beta)}. \end{aligned} \quad (43)$$

After utilizing [49, (06.05.03.0002.01) and (06.01.16.0006.01)], it is proved that holds

$$\frac{2^x \Gamma\left(\frac{x}{2}\right) \Gamma\left(\frac{x+1}{2}\right)}{\Gamma(x)} = 2\sqrt{\pi}, \quad (44)$$

so the final form of  $b_n^{\mu \rightarrow \infty}$  is derived as

$$b_n^{\mu \rightarrow \infty} = \frac{1}{\pi}. \quad (45)$$

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